



King Fahd University of Petroleum & Minerals

DEPARTMENT OF MATHEMATICAL SCIENCES

Technical Report Series

TR 083

April 1986

Distributional $(n+1)$ – Dimensional Heat Equation

M.A. Chaudhry and M.H. Kazi

DISTRIBUTIONAL $(n+1)$ — DIMENSIONAL HEAT EQUATION

M.A. Chaudhry and M.H. Kazi
 Department of Mathematical Sciences
 University of Petroleum & Minerals
 Dhahran, Saudi Arabia

ABSTRACT

Let $f(x) \in L_{\infty}(\mathbb{R})$. It is a classical result that the solution to the following initial value problem associated with the heat equation:

$$\begin{aligned} \frac{\partial U(x,t)}{\partial t} - c^2 \frac{\partial^2 U(x,t)}{\partial x^2} &= 0 \\ U(x,0^+) &= f(x) \end{aligned} \quad (1)$$

in the domain $\mathbb{R}^{2,+} = \{(x,t): x \in \mathbb{R}, t > 0\}$ is:

$$U(x,t) = \frac{1}{2c\sqrt{\pi t}} \int_{-\infty}^{\infty} f(\xi) \exp\left(-\frac{|x-\xi|^2}{4c^2 t}\right) d\xi, \quad (2)$$

$U(x,t)$ being the temperature function. We extend the notion of the temperature function to the space $\mathbb{R}^{n+1,+} = \{(x,t): x \in \mathbb{R}^n, t > 0\}$. We call $U(x,t)$ a temperature function in an open region of $\mathbb{R}^{n+1,+}$, if it is infinitely differentiable at each point of the region and satisfies:

$$U_t(x,t) - c^2 \sum_{i=1}^n U_{x_i x_i}(x,t) = 0$$

In this paper we exploit the above definition of the temperature function to solve the Heat equation in $\mathbb{R}^{n+1,+}$, with a distributional initial condition.

As it turns out, our solution is quite constructive and its two dimensional case is an extension of the classical solution to the above-mentioned initial value problem.

1980 Mathematics Subject Classification:

Primary 46F12

Secondary 44A15

Keywords and Phrases:

Heat equation, Distributional boundary value problem,
generalized functions and distributions.

SCHWARTZ TESTING FUNCTION SPACE $\mathcal{D}_{L_p}(\mathbb{R}^n)$ ($1 \leq p < \infty$).

An infinitely differentiable function ϕ defined over \mathbb{R}^n is said to belong to the space $\mathcal{D}_{L_p}(\mathbb{R}^n)$ if:

$$D^\alpha \phi(x) = \frac{\partial^{|\alpha|} \phi(x)}{\partial x_1^{\alpha_1} \partial x_2^{\alpha_2} \dots \partial x_n^{\alpha_n}}$$

where $|\alpha| = \alpha_1 + \alpha_2 + \dots + \alpha_n$, belongs to $L_p(\mathbb{R}^n)$ for all $|\alpha| \geq 0$.

We introduce a sequence of semi-norms $\{\gamma_m\}_{m=0}^\infty$ on $\mathcal{D}_{L_p}(\mathbb{R}^n)$ as follows: For $\phi \in \mathcal{D}_{L_p}(\mathbb{R}^n)$ we define:

$$\gamma_{|\alpha|}(\phi) = \left(\int_{-\infty}^{\infty} |D^\alpha \phi(x)|^p dx \right)^{1/p} \quad (3)$$

where $|\alpha| = 0, 1, 2, 3, \dots$

Since γ_0 is a norm, the sequence of semi-norms $\{\gamma_m\}_{m=0}^\infty$ is separating. The topology of $\mathcal{D}_{L_p}(\mathbb{R}^n)$ is generated by the semi-norms in the usual manner [1, pp. 8-14]. We say that a sequence $\{\phi_\mu\}_{\mu=1}^\infty$ converges to ϕ in $\mathcal{D}_{L_p}(\mathbb{R}^n)$, if for each $m = 0, 1, 2, 3, \dots$

$$\gamma_m(\phi_\mu - \phi) \rightarrow 0 \quad \text{as } \mu \rightarrow \infty$$

The space $\mathcal{D}_{L_p}(\mathbb{R}^n)$ is a complete countably multinormal space [2, p. 87]:

We denote by $(\mathcal{D}_{L_p}(\mathbb{R}^n))'$ the space of all continuous linear functionals on $\mathcal{D}_{L_p}(\mathbb{R}^n)$. We state here the well known result regarding the structure formula for $f \in (\mathcal{D}_{L_p}(\mathbb{R}^n))'$ [4, p.201]. See also [2, pp.109-116].

THEOREM 1: Let $f \in (\mathcal{D}_{L_p}(\mathbb{R}^n))'$ ($1 < p < \infty$) then f is equal to a finite linear combination of the derivatives of functions in $L_q(\mathbb{R}^n)$, that is for each $f \in (\mathcal{D}_{L_p}(\mathbb{R}^n))'$

$$\langle f, \phi \rangle = \sum_{|\alpha|=0}^r (-1)^{|\alpha|} \int_{\mathbb{R}^n} D^\alpha \phi \cdot f_{|\alpha|}(x) dx; \quad \forall \phi \in \mathcal{D}_{L_p}(\mathbb{R}^n) \quad (4)$$

where $f_{|\alpha|}$'s are functions in $L_q(\mathbb{R}^n)$, $\frac{1}{p} + \frac{1}{q} = 1$

Let us consider the Weierstrass kernel for $\mathbb{R}^{n+1,+}$ given by:

$$w(x,t) = \left(\frac{1}{\sqrt{4\pi t c^2}}\right)^n \exp\left(-\frac{|x|^2}{4tc^2}\right). \quad (5)$$

The following properties of $w(x,t)$ are well known [3]:

$$i) \quad \int_{\mathbb{R}^n} w(x,t) dx = 1 \quad \text{for each } t > 0 \quad (6)$$

ii) Let $\phi \in L_q(\mathbb{R}^n)$ then $U_\phi(k,t)$ defined by:

$$U_\phi(x,t) = \int_{\mathbb{R}^n} \phi(x - \xi) w(\xi,t) d\xi \quad (7)$$

is a temperature function in $\mathbb{R}^{n+1,+}$, and

$$a) \quad \|U_\phi(x,t)\|_{L_q} \leq A_t \|\phi(\lambda)\|_{L_q}, \quad t > 0$$

$$b) \quad U_\phi(x,t) \text{ converges to } \phi(x) \text{ in the norm of } L_q(\mathbb{R}^n).$$

THEOREM 2: Let $f \in (\mathcal{D}_{L_p}(\mathbb{R}^n))'$ ($1 < p < \infty$). Then, for all $\phi \in \mathcal{D}_{L_p}(\mathbb{R}^n)$,

$$\langle \langle f(\xi), w(x-\xi,t) \rangle, \phi(x) \rangle \rightarrow \langle f(x), \phi(x) \rangle \quad \text{as } t \rightarrow 0^+$$

i.e.

$$\langle f(\xi), w(x-\xi,t) \rangle \rightarrow f(x) \text{ in } (\mathcal{D}_{L_p}(\mathbb{R}^n))' \text{ as } t \rightarrow 0^+$$

PROOF: Using the structure formula for $f \in (\mathcal{D}_{L_p}(\mathbb{R}^n))'$ we have

$$\begin{aligned} I(t) &= \langle \langle f(\xi), w(x-\xi,t) \rangle, \phi(x) \rangle \\ &= \sum_{|\alpha|=0}^r (-1)^{|\alpha|} \langle f_{|\alpha|}(\xi), D_\xi^\alpha w(x-\xi,t) \rangle, \phi(x) \rangle \end{aligned}$$

where $f_{|\alpha|}$'s are in $L_q(\mathbb{R}^n)$.

However,

$$\frac{\partial}{\partial \xi_j} W(x-\xi,t) = - \frac{\partial}{\partial x_j} W(x-\xi,t)$$

Or more generally,

$$D_{\xi}^{\alpha} w(x-\xi, t) = (-1)^{|\alpha|} D_x^{\alpha} w(x-\xi, t). \quad (9)$$

Therefore, we can write (8) as:

$$I(t) = \sum_{|\alpha|=0}^r \langle \langle f_{|\alpha|}(\xi), D_x^{\alpha} w(x-\xi, t) \rangle, \phi(x) \rangle$$

Using Fubini's theorem and integration by parts, we get:

$$\begin{aligned} I(t) &= \sum_{|\alpha|=0}^r (-1)^{|\alpha|} \langle f_{|\alpha|}(\xi), \langle w(x-\xi, t), \phi^{|\alpha|}(x) \rangle \rangle \\ &= \sum_{|\alpha|=0}^r (-1)^{|\alpha|} \langle f_{|\alpha|}(\xi), \int_{\mathbb{R}^n} w(x-\xi, t) \cdot \phi^{|\alpha|}(x) \cdot dx \rangle \quad (10) \end{aligned}$$

$$\text{Let } U_{\phi, \alpha}(\xi, t) = \int_{\mathbb{R}^n} w(x-\xi, t) \phi^{(\alpha)}(x) dx$$

Then, it follows from (7) that for each fixed $t > 0$, $U_{\phi, \alpha}(\xi, t)$ belongs to $C^{\infty}(\mathbb{R}^{n+1}, +)$ and

$$\|U_{\phi, \alpha}(\xi, t)\|_{L_p} \leq C_{\phi} \|\phi^{(\alpha)}(\xi)\|_{L_p} < \infty \quad t > 0 \quad (11)$$

and

$$\|U_{\phi, \alpha}(\xi, t) - \phi^{(\alpha)}(\xi)\|_{L_p} \rightarrow 0 \quad \text{as } t \rightarrow 0^+. \quad (12)$$

However,

$$D^{\beta}[U_{\phi, \alpha}(\xi, t)] = \int_{\mathbb{R}^n} w(x-\xi, t) \cdot \phi^{(\alpha+\beta)}(x) dx$$

$$(\forall |\beta| = 0, 1, 2, \dots)$$

Therefore,

$$\gamma_{|\beta|} (U_{\phi, \alpha}(\xi, t) - \phi^{(\alpha)}(\xi)) \rightarrow 0 \quad \text{as } t \rightarrow 0 \quad (13)$$

$$(\forall |\beta| = 0, 1, 2, 3, \dots)$$

Thus,

$$U_{\phi, \alpha}(\xi, t) \xrightarrow{\mathcal{D}_{L_p}(\mathbb{R}^n)} \phi^{(\alpha)}(\xi) \quad \text{as } t \rightarrow 0. \quad (14)$$

From (10) and (14) we get

$$\begin{aligned} \lim_{t \rightarrow 0^+} I(t) &= \sum_{|\alpha|=0}^r (-1)^{|\alpha|} \langle f_{|\alpha|}(\xi), \phi^{(\alpha)}(\xi) \rangle \\ &= \langle f, \phi \rangle \quad \forall \phi \in \mathcal{D}_{L_p}(\mathbb{R}^n) \quad \text{as } t \rightarrow 0^+ \quad [\text{by 4}] \end{aligned}$$

EXAMPLE 1: Consider the distributional heat problem:

$$U_t(x, t) - c^2 \sum_{i=1}^n U_{x_i x_i}(x, t) = 0 \quad (15)$$

$$\lim_{t \rightarrow 0^+} U(x, t) = f(x) \quad \text{in } (\mathcal{D}_{L_p}(\mathbb{R}^n))' \quad (16)$$

Let us define a function $U(x, t)$ by:

$$U(x, t) = \langle f(\xi), w(x-\xi, t) \rangle, \quad t > 0$$

Using the structure formula (4) for f we get:

$$\begin{aligned}
 U(x,t) &= \sum_{|\alpha|=0}^r (-1)^{|\alpha|} \langle f_{|\alpha|}(\xi), D_{\xi}^{\alpha} w(x-\xi, t) \rangle \\
 &= \sum_{|\alpha|=0}^r (-1)^{|\alpha|} \cdot U_{\alpha}(x,t) \quad (17)
 \end{aligned}$$

According to (7), each $U_{\alpha}(x,t)$ is a temperature function, therefore $U(x,t)$ is a temperature function. In view of the theorem (2), $U(x,t)$ satisfies the distributional-initial condition (15).

EXAMPLE 2: Consider the case $n = 1$. The distributional initial value problem in $\mathbb{R}^{2,+} = \{(x,t) : x,t \in \mathbb{R}, t > 0\}$ becomes

$$U(x,t) - c^2 U_{xx}(x,t) = 0$$

$$U(x,0^+) = f(x) \text{ in } (\mathcal{D}'_L(\mathbb{R}))'$$

The above problem is the heat conduction problem for a straight infinite bar of uniform cross section and homogeneous material with the initial temperature distribution in the bar is given by the generalized function $f(x)$. The solution of the problem follows from (17) in the particular case when $n = 1$:

$$U(x,t) = \langle f(\xi), \frac{1}{2c\sqrt{\pi t}} \exp\left(-\frac{|x-\xi|^2}{4c^2 t}\right) \rangle.$$

References

1. A. Friedman, Generalized Functions and Partial Differential Equations, Prentice-Hall, Englewood Cliffs, N.J., (1963).
2. I.M. Gelfand, and G.E. Shilov, Generalized Functions, Vol.2, Moscow.
3. G.O. Okikiolu, Special Integral Operators: Vol.I. Weierstrass Operators and Related Integrals. Okikiolu Scientific and Industrial Organization (1980).
4. L. Schwartz, Theories des Distributions, Vol. I, II, Hermann Paris, (1957), (1959).
5. A.H. Zemanian, Generalized Integral Transforms, Inter Science Publishers, New York, (1968).