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Dam Problem**

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Abstract

We consider a non steady-state fluid flow through a heterogeneous porous medium governed by a nonlinear Darcy law. Under a general condition on the permeability we prove a continuity result for the function characterizing the wet part of the dam.

1 Formulation of the Problem

A porous medium that we denote by Ω is supplied by several reservoirs of a fluid which infiltrates through Ω . We assume that Ω is a bounded locally Lipschitz domain of \mathbb{R}^n with boundary $\partial\Omega = \Gamma_1 \cup \Gamma_2$, where Γ_1 is the impervious part, Γ_2 is the part in contact with either air or fluid reservoirs. Let T be a positive number, $Q = \Omega \times (0, T)$ and ϕ a nonnegative Lipschitz function defined on \overline{Q} .

We are concerned with the problem of finding the pressure p and the saturation χ of the fluid. For convenience we introduce the following functions : $\psi = \phi + x_n$, $u = p + x_n$ and $g = 1 - \chi$. The flow inside Ω is incompressible and obeys to the following generalized Darcy law (see [10])

$$v = -\mathcal{A}(x, \nabla(p + x_n))$$

where \mathcal{A} is a function defined in $\Omega \times \mathbb{R}^N$.

Using the mass conservation law, Darcy's law, the boundary conditions and the initial data, we obtain the following strong formulation for our problem (see [4])

$$(SF) \begin{cases} i) & u \geq x_n, \quad 0 \leq g \leq 1, \quad g \cdot (u - x_n) = 0, & \text{in } Q \\ ii) & \operatorname{div}(\mathcal{A}(x, \nabla u) - g\mathcal{A}(x, e)) + g_t = 0, & \text{in } Q \\ iii) & u = \psi, & \text{on } \Sigma_2 \\ iv) & g(\cdot, 0) = g_0, & \text{in } \Omega \\ v) & (\mathcal{A}(x, \nabla u) - g\mathcal{A}(x, e)) \cdot \nu = 0, & \text{on } \Sigma_1 \\ vi) & (\mathcal{A}(x, \nabla u) - g\mathcal{A}(x, e)) \cdot \nu \leq 0, & \text{on } \Sigma_4 \end{cases}$$

where $g_0 \in L^\infty(\Omega)$, $e = (0, \dots, 0, 1) \in \mathbb{R}^n$,

- $\Sigma_1 = \Gamma_1 \times (0, T)$: the impervious part,
- $\Sigma_2 = \Gamma_2 \times (0, T)$: the pervious part,
- $\Sigma_3 = \Sigma_2 \cap \{\phi > 0\}$: the part covered by fluid,
- $\Sigma_4 = \Sigma_4 \cap \{\phi = 0\}$: the part where the fluid flows outside Ω .

For \mathcal{A} we assume the following with $q > 1$ and $0 < m \leq M < \infty$:

$$\begin{cases} i) & X \mapsto \mathcal{A}(X, \xi) \text{ is measurable } \forall \xi \in \mathbb{R}^2, \\ ii) & \xi \mapsto \mathcal{A}(X, \xi) \text{ is continuous for a.e. } X \in \Omega, \\ iii) & \text{for all } \xi \in \mathbb{R}^2 \text{ and for a.e. } X \in \Omega, \\ & \mathcal{A}(X, \xi) \cdot \xi \geq m|\xi|^q \quad \text{and} \quad |\mathcal{A}(X, \xi)| \leq M|\xi|^{q-1}, \\ iv) & \text{for all } \xi, \zeta \in \mathbb{R}^2 \text{ and for a.e. } X \in \Omega, \\ & (\mathcal{A}(X, \xi) - \mathcal{A}(X, \zeta)) \cdot (\xi - \zeta) \geq 0. \end{cases} \quad (1.1)$$

Using the strong formulation, we are led to the following weak formulation

$$(P) \begin{cases} \text{Find } (u, g) \in L^q(0, T, W^{1,q}(\Omega)) \times L^\infty(Q) \text{ such that :} \\ (i) & u \geq x_n, \quad 0 \leq g \leq 1, \quad g(u - x_n) = 0 \quad \text{a.e. in } Q, \\ (ii) & u = \psi \quad \text{on } \Sigma_2, \\ (iii) & \int_Q (\mathcal{A}(X, \nabla u) - g\mathcal{A}(X, e)) \cdot \nabla \xi + g\xi_t dx dt + \int_\Omega g_0(x)\xi(x, 0) dx \leq 0 \\ & \forall \xi \in W^{1,q}(Q), \quad \xi = 0 \text{ on } \Sigma_3, \quad \xi \geq 0 \text{ on } \Sigma_4, \quad \xi(x, T) = 0 \quad \text{a.e. in } \Omega. \end{cases}$$

Under the assumptions (1.1), the existence of a solution was proved in [11] and also in [4] for generalized boundary conditions. Here we are concerned with the

continuity of the function g . Assuming that $\mathcal{A}(X, e)$ is a constant vector, it was proved (see [11]) that $g \in C^0([0, T], L^p(\Omega))$ for all $p \geq 1$. The same result was proved in [3] and [2] in the case where $\mathcal{A}(X, \xi) = \xi$. Our main objective in this paper is to extend this regularity result to the case where $\operatorname{div}(\mathcal{A}(X, e)) \geq 0$. The idea of the proof is based on a monotonicity result of g along the orbits of a certain differential equation. A similar monotonicity is proved in [5] for the stationary case.

We recall the following results from [11]

Proposition 1.1. *For each solution (u, g) of (P), we have*

$$u \in L^\infty(Q) \quad \text{and} \quad g \in C^0([0, T], W^{-1, q'}(\Omega)), \quad (1.2)$$

$$\operatorname{div}(\mathcal{A}(X, \nabla u) - g\mathcal{A}(X, e)) + g_t = 0 \quad \text{in } \mathcal{D}'(Q). \quad (1.3)$$

Moreover if $\operatorname{div}(\mathcal{A}(X, e)) \geq 0$ in $\mathcal{D}'(\Omega)$, we obtain

$$\operatorname{div}(\mathcal{A}(X, \nabla u)) - g_t = \operatorname{div}(g\mathcal{A}(X, e)) \geq 0 \quad \text{in } \mathcal{D}'(Q). \quad (1.4)$$

2 A Monotonicity Property of g

From now on, we assume that

$$\mathcal{A}(\cdot, e) = (a^1(\cdot), a^2(\cdot)) \in C^1(\bar{\Omega}) \quad (2.1)$$

$$\operatorname{div}(\mathcal{A}(X, e)) \geq 0 \quad \text{in } C^0(\Omega) \quad (2.2)$$

$$\partial\Omega \quad \text{is of class } C^1 \quad (2.3)$$

$$\mathcal{A}(X, e) \cdot \nu \neq 0 \quad \forall X \in \partial\Omega. \quad (2.4)$$

Then we consider the following differential system

$$(E(\omega, h)) \quad \begin{cases} X'(s, \omega, h) &= \mathcal{A}(X(s, \omega, h), e) \\ X(0, \omega, h) &= (\omega, h) \end{cases}$$

where $h \in \pi_{x_n}(\Omega)$, $\omega \in \pi_{x'}(\Omega \cap [x_n = h])$ and where $\pi_{x'}$ and π_{x_n} are respectively the orthogonal projections on the hyperplane $x_n = 0$ and the x_n axis.

By the classical theory of ordinary differential equations there exists a unique maximal solution $X(\cdot, \omega, h)$ of $E(\omega, h)$ which is defined on $[\alpha_-(\omega, h), \alpha_+(\omega, h)]$ with

$$X(\alpha_-(\omega, h), \omega, h) \in \partial\Omega \cap ([x_n < h]) \quad \text{and} \quad X(\alpha_+(\omega, h), \omega, h) \in \partial\Omega \cap ([x_n > h]).$$

Under the assumptions (2.1), (2.3) and (2.4), one can prove (see [5]) that $\alpha_-, \alpha_+ \in C^1(\pi_{x'}(\Omega \cap [x_n = h]))$. This allows us to introduce the following definition as in [5].

Definition 2.1. For each $h \in \pi_{x_n}(\Omega)$ we define the set

$$D_h = \{(s, \omega) / \omega \in \pi_{x'}(\Omega \cap [x_n = h]), s \in (\alpha_-(\omega, h), \alpha_+(\omega, h))\}$$

and consider the mappings $T_h : D_h \longrightarrow T_h(D_h)$ and $S_h : D_h \longrightarrow S_h(D_h)$ defined by

$$T_h(s, \omega) = X(s, \omega, h) \quad \text{and} \quad S_h(s, \omega) = (\omega, L_h(s, \omega)) = (\omega, \tau),$$

where $L_h(s, \omega) = \int_{\alpha_-(\omega, h)}^s |\mathcal{A}(X(\nu, \omega, h), e)| d\nu = \int_{\alpha_-(\omega, h)}^s |X'(\nu, \omega, h)| d\nu$ represents the arc Length of the curve $X(\cdot, \omega, h)$ from the point $X(\alpha_-(\omega, h), \omega, h)$ to the point $X(s, \omega, h)$.

Then we have

Proposition 2.1.

$$\Omega = \bigsqcup_{h \in \pi_{x_n}(\Omega)} T_h(D_h), \quad T_h \quad \text{and} \quad S_h \quad \text{are} \quad C^1 \quad \text{diffeomorphisms.}$$

Proof. We refer to [5] for the proof and we recall the following properties

$$\det \mathcal{J}S_h = -|\mathcal{A}(X(t, \omega, h), e)| < 0$$

$$Y_h(t, \omega) = \det(\mathcal{J}T_h) = -a^2(X(0, \omega, h)) \cdot \exp\left(\int_0^t \{div(\mathcal{A}(\cdot, e))\}(X(s, \omega, h)) ds\right) < 0,$$

where $\mathcal{J}F$ is the jacobian of F . □

The following monotonicity result generalizes the fact that $g_{x_n} - g_t \geq 0$ in $\mathcal{D}'(Q)$ when $\mathcal{A}(X, \xi) = \xi$ (see [2]). It will play a major role for the continuity proof of g .

Theorem 2.1. Let (u, g) be a solution of (P). We have for each $h \in \pi_{x_n}(\Omega)$

$$f_\tau - \lambda f_t \geq 0 \quad \text{in} \quad \mathcal{D}'(S_h(D_h) \times (0, T))$$

where $f = goT_hoS_h^{-1} \cdot (-Y_hoS_h^{-1})$ and $\lambda(\omega, \tau) = |\mathcal{A}(T_hoS_h^{-1}(\omega, \tau), e)|^{-1}$.

Proof. Let $\phi \in \mathcal{D}(S_h(D_h) \times (0, T))$, $\phi \geq 0$. Then $\tilde{\phi}(x', x_n, t) = \phi(S_h \circ T_h^{-1}(x', x_n), t) \in C_0^1(T_h(D_h) \times (0, T))$, $\tilde{\phi} \geq 0$ and by (1.4) and (2.2), we have

$$\int_{T_h(D_h) \times (0, T)} g \mathcal{A}(X, e) \cdot \nabla \tilde{\phi} - g \tilde{\phi}_t dx dt \leq 0$$

which can be written

$$\int_{T_h(D_h) \times (0, T)} g \mathcal{A}(X, e) \cdot \nabla (\phi \circ S_h \circ T_h^{-1}) - g \phi_t (S_h \circ T_h^{-1}(x', x_n), t) dx' dx_n dt \leq 0.$$

Using the change of variables $(T_h(s, \omega), t) = (x', x_n, t)$ and the fact that

$$\mathcal{A}(X(s, \omega), e) (\nabla (\phi \circ S_h \circ T_h^{-1})) \circ T_h \cdot (-Y_h(s, \omega)) = -Y_h(s, \omega) \frac{\partial}{\partial s} (\phi \circ S_h)$$

we get

$$\begin{aligned} & \int_{D_h \times (0, T)} g(T_h(s, \omega), t) \cdot (-Y_h(s, \omega)) \cdot \frac{\partial}{\partial s} \phi(S_h(s, \omega), t) ds d\omega dt \\ & - \int_{D_h \times (0, T)} g(T_h(s, \omega), t) \cdot (-Y_h(s, \omega)) \cdot \frac{\partial}{\partial t} \phi(S_h(s, \omega), t) ds d\omega dt \leq 0 \end{aligned}$$

which becomes after using the change of variables S_h^{-1}

$$\begin{aligned} & \int_{S_h(D_h) \times (0, T)} g(T_h \circ S_h^{-1}(\omega, \tau), t) \cdot (-Y_h \circ S_h^{-1}(\omega, \tau)) \cdot \left(\frac{\partial}{\partial s} (\phi \circ S_h(\cdot, \cdot)) \right) (S_h^{-1}(\omega, \tau), t) \\ & \quad \cdot |det \mathcal{J} S_h^{-1}(\omega, \tau)| d\omega d\tau dt \\ & - \int_{S_h(D_h) \times (0, T)} g(T_h \circ S_h^{-1}(\omega, \tau), t) \cdot (-Y_h \circ S_h^{-1}(\omega, \tau)) \cdot |det \mathcal{J} S_h^{-1}(\omega, \tau)| \\ & \quad \cdot \frac{\partial \phi}{\partial t}(\omega, \tau, t) \cdot d\omega d\tau dt \leq 0. \end{aligned}$$

Taking into account the fact that

$$\left(\frac{\partial}{\partial s} (\phi \circ S_h(\cdot, \cdot)) \right) (S_h^{-1}(\omega, \tau), t) = \frac{\partial \phi}{\partial \tau} \cdot \mathcal{A}(\cdot, e) \circ T_h \circ S_h^{-1}(\omega, \tau) = \frac{\partial \phi}{\partial \tau} \cdot |det \mathcal{J} S_h|,$$

we obtain

$$\int_{S_h(D_h) \times (0, T)} f \cdot \frac{\partial \phi}{\partial \tau} - |\mathcal{A}(T_h \circ S_h^{-1}(\omega, \tau), e)|^{-1} f \cdot \frac{\partial \phi}{\partial t} d\omega d\tau dt \leq 0.$$

□

As a consequence of Theorem 1.1, we have the following monotonicity result

Theorem 2.2. *Let (u, g) be a solution of (P). For each $h \in \pi_{x_n}(\Omega)$ and $\xi \in \mathcal{D}(S_h(D_h) \times (0, T))$, $\xi \geq 0$, the function*

$$F(k) = \int_{\text{supp}(\xi)} f\left(\omega, \tau + k, t - \int_{\tau}^{\tau+k} \lambda(\omega, \sigma) d\sigma\right) \xi(\omega, \tau, t) d\omega d\tau dt$$

is non-increasing on its domain of definition.

Proof. First of all we have with $Q_h = S_h(D_h) \times (0, T)$

$$F(k) = \int_{\vartheta_k^{-1}(Q_h)} f\left(\omega, \tau + k, t - \int_{\tau}^{\tau+k} \lambda(\omega, \sigma) d\sigma\right) \xi(\omega, \tau, t) d\omega d\tau dt \quad (2.5)$$

where $\vartheta_k(\omega, \tau, t) = \left(\omega, \tau + k, t - \int_{\tau}^{\tau+k} \lambda(\omega, \sigma) d\sigma\right)$. Moreover ϑ_k is differentiable with $\mathcal{J}\vartheta_k(\omega, \tau, t) = 1$. Therefore we obtain from (2.5) by using the change of variables $\vartheta_k(\omega, \tau, t) = (\omega, \nu, s)$

$$F(k) = \int_{Q_h} f(\omega, \nu, s) \xi\left(\omega, \nu - k, s + \int_{\nu-k}^{\nu} \lambda(\omega, \sigma) d\sigma\right) d\omega d\nu ds. \quad (2.6)$$

From (2.6) we deduce that F is differentiable with

$$\begin{aligned} F'(k) &= \int_{Q_h} f(\omega, \nu, s) \left\{ -\frac{\partial \xi}{\partial \tau} + \lambda(\omega, \nu - k) \frac{\partial \xi}{\partial t} \right\} \left(\omega, \nu - k, s + \int_{\nu-k}^{\nu} \lambda(\omega, \sigma) d\sigma \right) d\omega d\nu ds \\ &= \int_{Q_h} f(\omega, \nu, s) \left\{ -\frac{\partial \zeta}{\partial \nu} + \lambda(\omega, \nu - k) \frac{\partial \zeta}{\partial s} \right\} (\omega, \nu, s) d\omega d\nu ds, \end{aligned}$$

where $\zeta(\omega, \nu, s) = \xi\left(\omega, \nu - k, s + \int_{\nu-k}^{\nu} \lambda(\omega, \sigma) d\sigma\right)$.

Since $\zeta \in \mathcal{C}_0^1(Q_h)$ and $\zeta \geq 0$, it follows from Theorem 2.1 that $F'(k) \leq 0$. □

3 Continuity of g

The main result is the following theorem

Theorem 3.1. *Let (u, g) be a solution of (P) . We have $g \in C^0([0, T], L^p(\Omega))$ for all $p \geq 1$.*

The proof of Theorem 3.1 is based on Theorem 2.2 and the following lemma.

Lemma 3.1. *Let (u, g) be a solution of (P) . We have $f \in C^0([0, T], L^2(S_h(D_h)))$.*

Proof. First of all we deduce from Proposition 1.1 and the fact that g is bounded

$$f(\omega, \tau, t + \epsilon) \xrightarrow{\epsilon \rightarrow 0} f(\omega, \tau, t) \text{ weakly in } L^2(S_h(D_h)). \quad (3.1)$$

Let $K = S_h(D_h)$, $\Lambda(\omega, \tau) = \int_{\tau_0}^{\tau} \lambda(\omega, \sigma) d\sigma$, (τ_0 fixed) and consider for $\epsilon > 0$ small enough

$$K_\epsilon = \{(\omega, \tau) \in \Omega / (\omega, \tau + \eta(\omega, \tau)) \in K\}$$

where η is defined by $\epsilon = \Lambda(\omega, \tau + \eta(\omega, \tau)) - \Lambda(\omega, \tau) = \int_{\tau}^{\tau + \eta} \lambda(\omega, \sigma) d\sigma$.

Note that it is enough to show that

$$\lim_{\epsilon \rightarrow 0} \int_K f^2(\omega, \tau, t + \epsilon) - f^2(\omega, \tau, t) d\omega d\tau. \quad (3.2)$$

To do this we first remark that

$$\begin{aligned} & \left| \int_K f^2(\omega, \tau, t + \epsilon) - f^2(\omega, \tau, t) d\omega d\tau \right| \leq \left| \int_{K \cap K_\epsilon} f^2(\omega, \tau, t + \epsilon) - f^2(\omega, \tau, t) d\omega d\tau \right| \\ & + \left| \int_{(K \setminus K_\epsilon) \cup (K_\epsilon \setminus K)} f^2(\omega, \tau, t + \epsilon) - f^2(\omega, \tau, t) d\omega d\tau \right| = I_{\epsilon 1} + I_{\epsilon 2}. \end{aligned} \quad (3.3)$$

Note that $I_{\epsilon 2} \leq C(|K \setminus K_\epsilon| + |K_\epsilon \setminus K|)$ (here and after we denote by C any positive constant) and therefore

$$\lim_{\epsilon \rightarrow 0} I_{\epsilon 2} = 0. \quad (3.4)$$

Moreover

$$\begin{aligned}
I_{\epsilon 1} &\leq \left| \int_{K \cap K_\epsilon} f^2(\omega, \tau, t + \epsilon) - f^2(\omega, \tau + \eta, t) d\omega d\tau \right| \\
&\quad + \left| \int_{K \cap K_\epsilon} f^2(\omega, \tau + \eta, t) - f^2(\omega, \tau, t) d\omega d\tau \right| \\
&= I_{\epsilon 3} + I_{\epsilon 4}
\end{aligned} \tag{3.5}$$

Note that by Theorem 2.2 we have

$$f(\omega, \tau, t + \epsilon) \geq f(\omega, \tau + \eta, t) \text{ a.e. } (\omega, \tau, t) \in Q_h.$$

Therefore, after using the fact that for a.e. $(\omega, \tau, t) \in Q_h$

$$|f^2(\omega, \tau, t + \epsilon) - f^2(\omega, \tau + \eta, t)| \leq C|f(\omega, \tau, t + \epsilon) - f(\omega, \tau + \eta, t)|,$$

we get

$$\begin{aligned}
I_{\epsilon 3} &\leq C \int_{K \cap K_\epsilon} |f(\omega, \tau, t + \epsilon) - f(\omega, \tau + \eta, t)| d\omega d\tau \\
&= C \int_{K \cap K_\epsilon} (f(\omega, \tau, t + \epsilon) - f(\omega, \tau + \eta, t)) d\omega d\tau \\
&= C \int_{K \cap K_\epsilon} (f(\omega, \tau, t + \epsilon) - f(\omega, \tau, t)) d\omega d\tau \\
&\quad + C \int_{K \cap K_\epsilon} (f(\omega, \tau, t) - f(\omega, \tau + \eta, t)) d\omega d\tau = CI_{\epsilon 5} + CI_{\epsilon 6}
\end{aligned} \tag{3.6}$$

Now

$$\begin{aligned}
|I_{\epsilon 5}| &\leq \int_K (f(\omega, \tau, t + \epsilon) - f(\omega, \tau, t)) d\omega d\tau \\
&\quad - \int_{K \setminus K_\epsilon} (f(\omega, \tau, t + \epsilon) - f(\omega, \tau, t)) d\omega d\tau \\
&\leq \int_K (f(\omega, \tau, t + \epsilon) - f(\omega, \tau, t)) d\omega d\tau + C|K \setminus K_\epsilon|.
\end{aligned} \tag{3.7}$$

It follows from (3.1) and (3.7) that

$$\lim_{\epsilon \rightarrow 0} I_{\epsilon 5} = 0. \quad (3.8)$$

Regarding $I_{\epsilon 6}$, we obtain by using the change of variables $(\omega, \tau) \rightarrow G(\omega, \tau) = (\omega, \tau + \eta(\omega, \tau)) = (\omega', \tau')$

$$I_{\epsilon 6} = \int_{K \cap K_\epsilon} f(\omega, \tau, t) d\omega d\tau - \int_{G(K \cap K_\epsilon)} f(\omega', \tau', t) \frac{\lambda(\omega', \tau')}{\lambda(\omega', \tau' - \eta(\omega, \tau))} d\omega' d\tau'.$$

Therefore it is clear that

$$\lim_{\epsilon \rightarrow 0} I_{\epsilon 6} = 0. \quad (3.9)$$

In the same way we prove that

$$\lim_{\epsilon \rightarrow 0} I_{\epsilon 4} = 0. \quad (3.10)$$

Taking into account (3.3)-(3.10), we get (3.2) and therefore the lemma is proved. \square

Proof of Theorem 3.1. Since $-Y_h \circ S_h^{-1}$ is positive, uniformly bounded and independent of t , we deduce from Lemma 3.1 that $g \circ T_h \circ S_h^{-1} = \frac{f}{(-Y_h \circ S_h^{-1})} \in C^0([0, T], L^2(S_h(D_h)))$. Moreover by using the change of variables $T_h \circ S_h^{-1}$ it follows that

$$g \in C^0([0, T], L^2(T_h(D_h))). \quad (3.11)$$

Now let $(\Omega_k)_{k \geq 1}$ be a monotone covering of Ω with open subsets satisfying $\overline{\Omega_k} \subset \Omega$. Since $(T_h(D_h))_{h \in \pi_{x_n}(\Omega)}$ is an open covering of Ω , we can cover each Ω_k by a finite number of sets $T_{h_i}(D_{h_i})$ and then it is clear that

$$g \in C^0([0, T], L^2(\Omega_k)), \quad \forall k \geq 1. \quad (3.12)$$

Now let $\epsilon > 0$ small enough and set $v_\epsilon(\omega, \tau, t) = |g(\omega, \tau, t + \epsilon) - g(\omega, \tau, t)|^2$. Let $\Omega'_1 = \Omega_1$, $\Omega'_{k+1} = \Omega_{k+1} \setminus \Omega_k$ for $k \geq 1$. Then $(\Omega'_k)_{k \geq 1}$ is a partition of Ω and we have

$$\int_{\Omega} v_\epsilon(\omega, \tau, t) d\omega d\tau = \sum_{k=1}^{\infty} \int_{\Omega'_k} v_\epsilon(\omega, \tau, t) d\omega d\tau. \quad (3.13)$$

Since $0 \leq \int_{\Omega'_k} v_\epsilon(\omega, \tau, t) d\omega d\tau \leq C|\Omega'_k|$ and $\sum_{k=1}^{\infty} C|\Omega'_k| = C|\Omega|$, it follows that

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega} v_{\epsilon}(\omega, \tau, t) d\omega d\tau = \sum_{k=1}^{\infty} \lim_{\epsilon \rightarrow 0} \int_{\Omega'_k} v_{\epsilon}(\omega, \tau, t) d\omega d\tau = 0. \quad (3.14)$$

Thus $g \in C^0([0, T], L^2(\Omega))$. Using the imbedding $L^2(\Omega) \subset L^p(\Omega)$, we obtain $g \in C^0([0, T], L^p(\Omega))$ for $p \in [1, 2]$.

Now for $p > 2$, we obtain the result since we have

$$\begin{aligned} |g(\omega, \tau, t + \epsilon) - g(\omega, \tau, t)|^p &= |g(\omega, \tau, t + \epsilon) - g(\omega, \tau, t)|^{p-2} |g(\omega, \tau, t + \epsilon) - g(\omega, \tau, t)|^2 \\ &\leq C |g(\omega, \tau, t + \epsilon) - g(\omega, \tau, t)|^2. \end{aligned}$$

□

Remark 3.1. *All results of this paper are clearly valid for the evolution dam problem with leaky boundary conditions ([12]).*

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